



Rapid Communication

# Chaos synchronization of double Duffing systems with parameters excited by a chaotic signal

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## Abstract

Chaos synchronization by driving parameter for two uncoupled identical chaotic double Duffing systems is presented. Replacing two corresponding parameters of the identical systems by the same function of chaotic state variables of a third chaotic system, the synchronization or anti-synchronization of two uncoupled systems can be obtained. Numerical simulations are illustrated for either synchronization or anti-synchronization of which the occurrence depends significantly on initial conditions and on driving strength. Alternative complete synchronization and anti-synchronization is also discovered.

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## 1. Introduction

Various synchronization phenomena are being reported for chaotic systems, such as complete synchronization (CS), anti-synchronization (AS), phase synchronization (PS), lag synchronization, and generalized synchronization [1–20,29–38]. However, most of synchronizations can only realize under the condition that there exists coupling between two chaotic systems. In practice, such as in physical and electrical systems, sometimes, it is difficult even impossible to couple two chaotic systems. In comparison with coupled chaotic systems, synchronization between the uncoupled chaotic systems has many advantages [20–29].

In this paper, synchronization of two double Duffing systems whose corresponding parameter is driven by a chaotic signal of a third system is analyzed. The chaos synchronizations of two uncoupled double Duffing systems are obtained by replacing their corresponding parameters by the same function of chaotic state variables of a third chaotic system. It is noted that whether CS or AS appear depends on the initial conditions. Besides, CS and AS are also characterized by great sensitivity to initial conditions and on the strengths of the substituted variable. It is found that CS or AS alternatively occurs under certain conditions [38–42].

This paper is organized as follows. In Section 2, a brief description of synchronization scheme based on the substitution of the strengths of the mutual coupling term of two identical chaotic double Duffing systems by

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the variable of a third double Duffing system is presented. In Section 3, numerical simulations are given for illustration. It is found that one can obtain CS or AS by adjusting the driving strength and initial conditions. Finally, in Section 4 conclusions are drawn.

## 2. Synchronization of two double Duffing systems

The famous Duffing system is

$$\ddot{x} + a\dot{x} + bx + cx^3 = d \cos \omega t \quad (1)$$

where  $a$ ,  $b$  are constant parameters,  $d \cos \omega t$  is an external excitation. It can be written as two first-order differential equations:

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = -ay - bx - cx^3 + d \cos \omega t \end{cases} \quad (2)$$

Consider the following double Duffing system:

$$\begin{cases} \frac{dx}{dt} = y \\ \frac{dy}{dt} = -ay - bx - cx^3 + du \\ \frac{du}{dt} = v \\ \frac{dv}{dt} = -ev - gu - hu^3 + kx \end{cases} \quad (3)$$

It consists of two Duffing systems in which two external excitations are replaced by two coupling terms. It is an autonomous system with four states where  $a$ ,  $b$ ,  $c$ ,  $d$ ,  $e$ ,  $g$ ,  $h$ , and  $k$  are constant parameters of the systems. Two identical double Duffing systems to be synchronized are

$$\begin{cases} \frac{dx_1}{dt} = y_1 \\ \frac{dy_1}{dt} = -ay_1 - bx_1 - cx_1^3 + d_1u_1 \\ \frac{du_1}{dt} = v_1 \\ \frac{dv_1}{dt} = -ev_1 - gu_1 - hu_1^3 + kx_1 \end{cases} \quad (4)$$

$$\begin{cases} \frac{dx_2}{dt} = y_2 \\ \frac{dy_2}{dt} = -ay_2 - bx_2 - cx_2^3 + d_2u_2 \\ \frac{du_2}{dt} = v_2 \\ \frac{dv_2}{dt} = -ev_2 - gu_2 - hu_2^3 + kx_2 \end{cases} \quad (5)$$

where  $a, b, c, d, e, g, h,$  and  $k$  are positive scalars, and  $d_1 = d_2$  are the control inputs to be designed. The third system is also a double Duffing system:

$$\begin{cases} \frac{dx_3}{dt} = y_3 \\ \frac{dy_3}{dt} = -ay_3 - bx_3 - cx_3^3 + du_3 \\ \frac{du_3}{dt} = v_3 \\ \frac{dv_3}{dt} = -ev_3 - gu_3 - hu_3^3 + kx_3 \end{cases} \quad (6)$$

In order to obtain the chaos synchronization of systems (4) and (5), the corresponding parameters  $d_1 = d_2$  of two systems are replaced by a chaotic signal  $px_3 + qv_3$  of the third system (6), where  $p, q$  are constant driving strengths. The error state variables are defined:

$$\begin{cases} e_1 = x_1 - x_2 \\ e_2 = y_1 - y_2 \\ e_3 = u_1 - u_2 \\ e_4 = v_1 - v_2 \end{cases} \quad (7)$$

Giving suitable values for  $p, q$  and initial conditions, the synchronization or anti-synchronization of systems (4) and (5) can be obtained.

### 3. Numerical simulations

Matlab method is used to all of our simulations with time step 0.01. The parameters of two systems (4) and (5) are given as  $a = 0.5, b = 1, c = 3, d = -2, e = 5, g = 1, h = 2, k = 2$  to ensure the chaotic behavior. To verify CS and AS, it is convenient to introduce the following coordinate transformation:  $E_1 = (x_1 + x_2)$  and  $e_1 = (x_1 - x_2)$  and the same transformation for  $y, u,$  and  $v$  variables. Therefore, the new coordinate systems  $(E_1, E_2, E_3, E_4)$  and  $(e_1, e_2, e_3, e_4)$  represent the sum and difference motions of the original coordinate system,

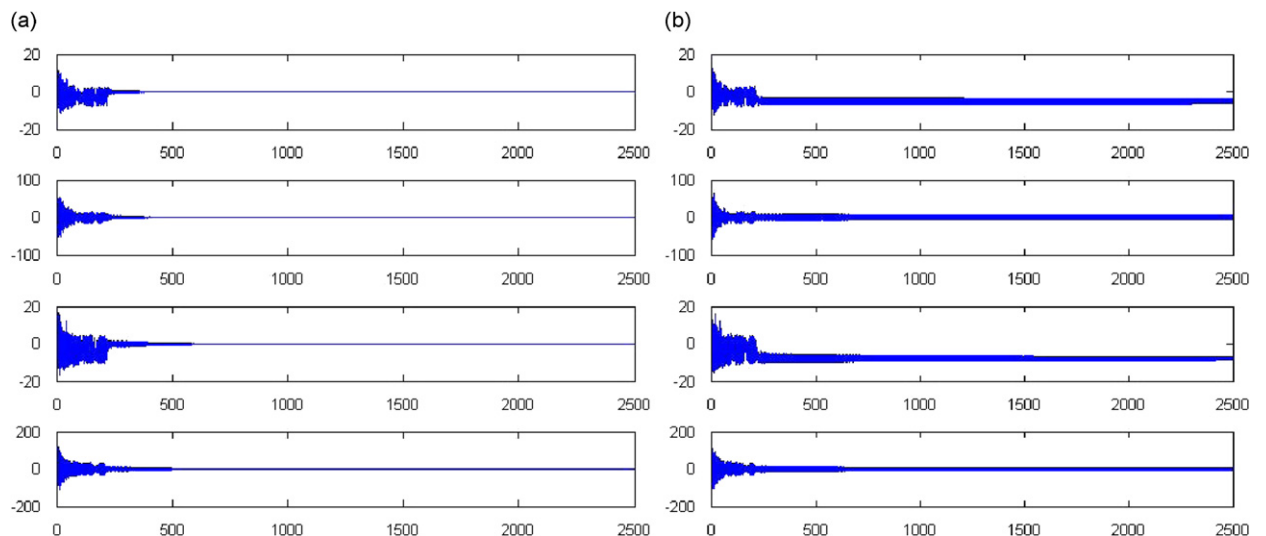


Fig. 1. CS and AS for initial condition  $(x_2, y_2, u_2, v_2) = (-8, -9, 0, 5)$ , and  $p = 10, q = 8$ . (a)  $e_1, e_2, e_3, e_4$  and (b)  $E_1, E_2, E_3, E_4$ .

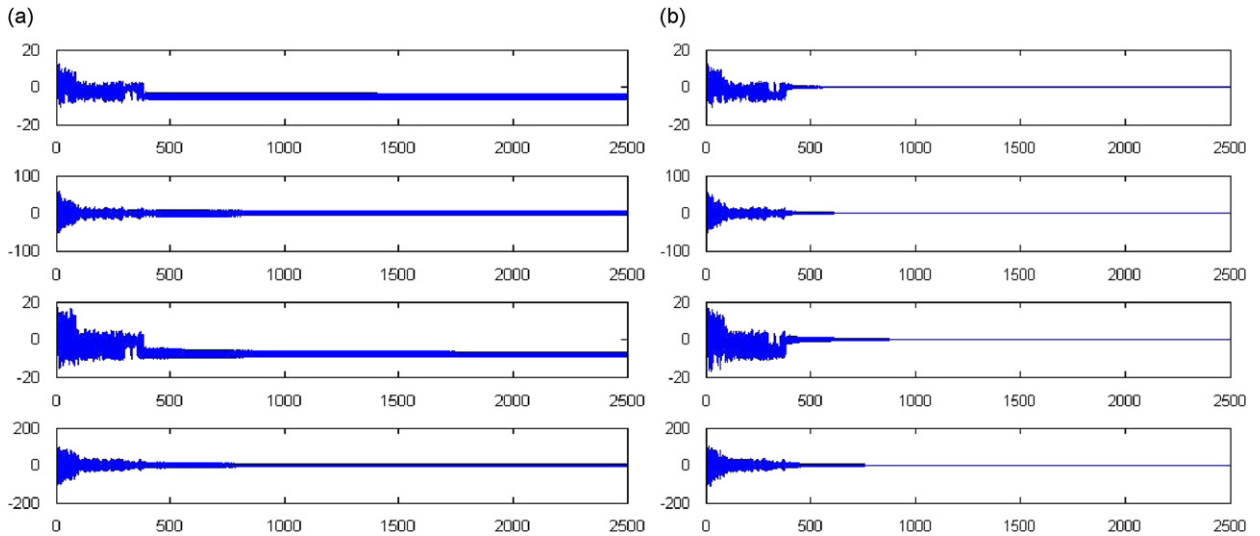


Fig. 2. CS and AS for initial condition  $(x_2, y_2, u_2, v_2) = (-8, -9, 0, 5)$ , and  $p = 10, q = 10$ . (a)  $e_1, e_2, e_3, e_4$  and (b)  $E_1, E_2, E_3, E_4$ .

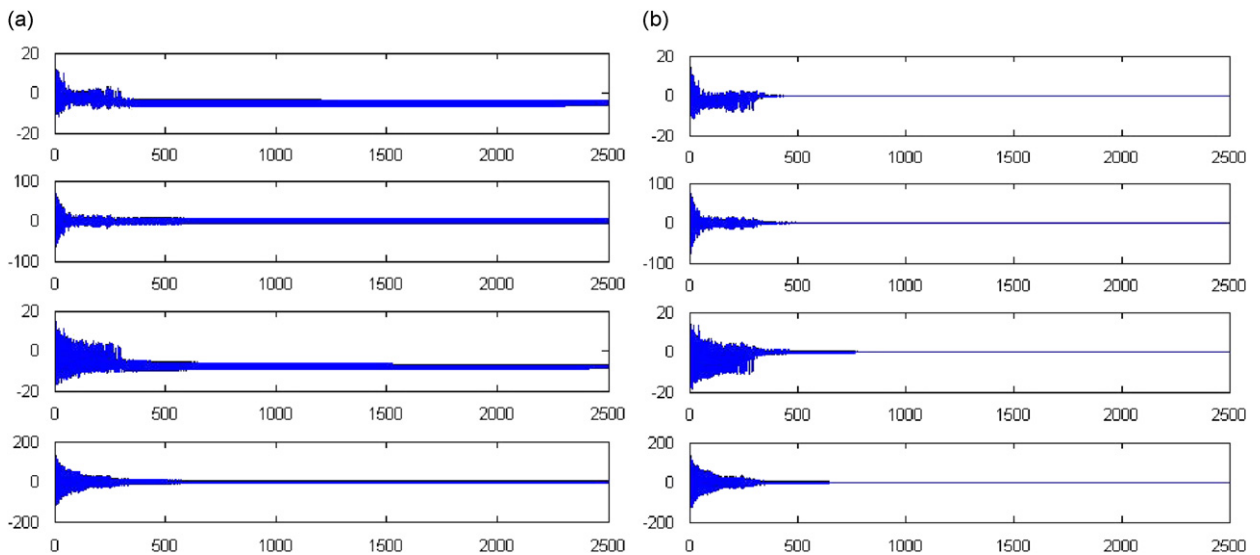


Fig. 3. CS and AS for initial condition  $(x_2, y_2, u_2, v_2) = (9, 5, -7, 9)$ , and  $p = 10, q = 10$ . (a)  $e_1, e_2, e_3, e_4$  and (b)  $E_1, E_2, E_3, E_4$ .

respectively. We can easily see that  $(e_1, e_2, e_3, e_4)$  subspace represents the CS case, and  $(E_1, E_2, E_3, E_4)$  subspace the AS one.

How the synchronization phenomena depend on the initial conditions will be studied. At the beginning, we choose  $(x_1, y_1, u_1, v_1) = (2, 5, 1, 0.3)$  and  $(x_2, y_2, u_2, v_2) = (-8, -9, 0, 5)$  as the initial conditions of systems (4) and (5). Let the driving strengths be  $p = 10, q = 8$  and  $p = 10, q = 10$ . Figs. 1 and 2 show the time-series of AS and CS phenomena for different driving strengths, respectively. The simulation results are shown in Fig. 1 for case (a) and in Fig. 2 for case (b). These simulation results indicate that the final state develops to CS or AS depending sensitively on driving strength in spite of the identical initial conditions in both cases. For AS case (Figs. 1(a) and (b)), the sums of the variables converge to zero, while the differences remain chaotic. For CS case (Figs. 2(a) and (b)), on the other hand,  $e_1, e_2, e_3,$  and  $e_4$  converge to zero, while  $E_1, E_2, E_3,$  and  $E_4$  remain chaotic.

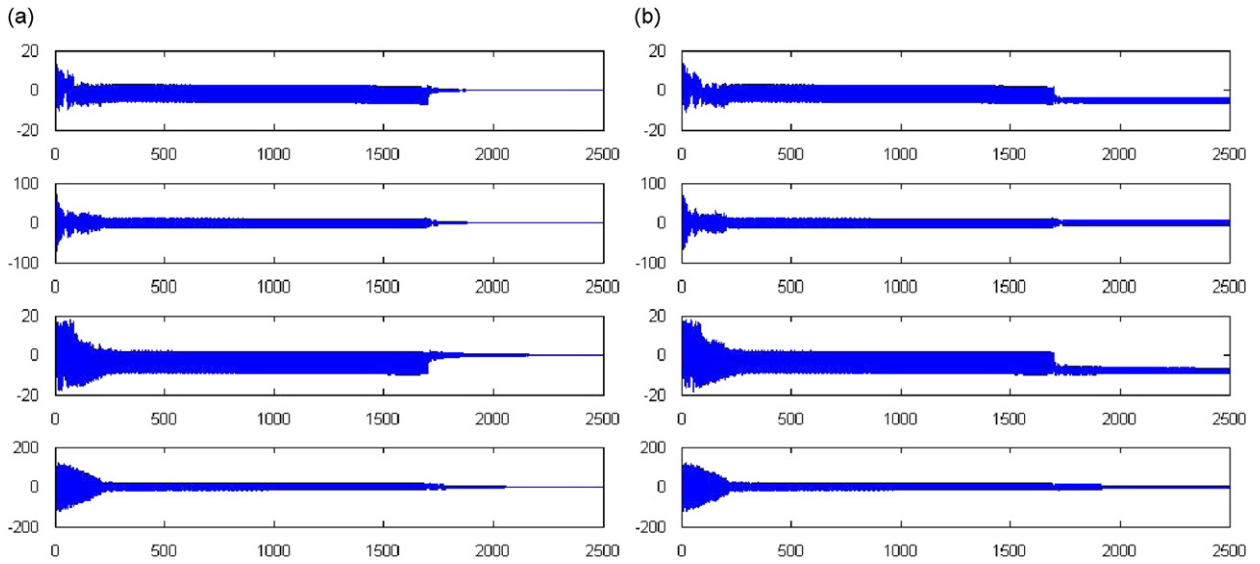


Fig. 4. CS and AS for initial condition  $(x_2, y_2, u_2, v_2) = (-8, -9, 0, 5)$ , and  $p = 10, q = 13$ . (a)  $e_1, e_2, e_3, e_4$  and (b)  $E_1, E_2, E_3, E_4$ .

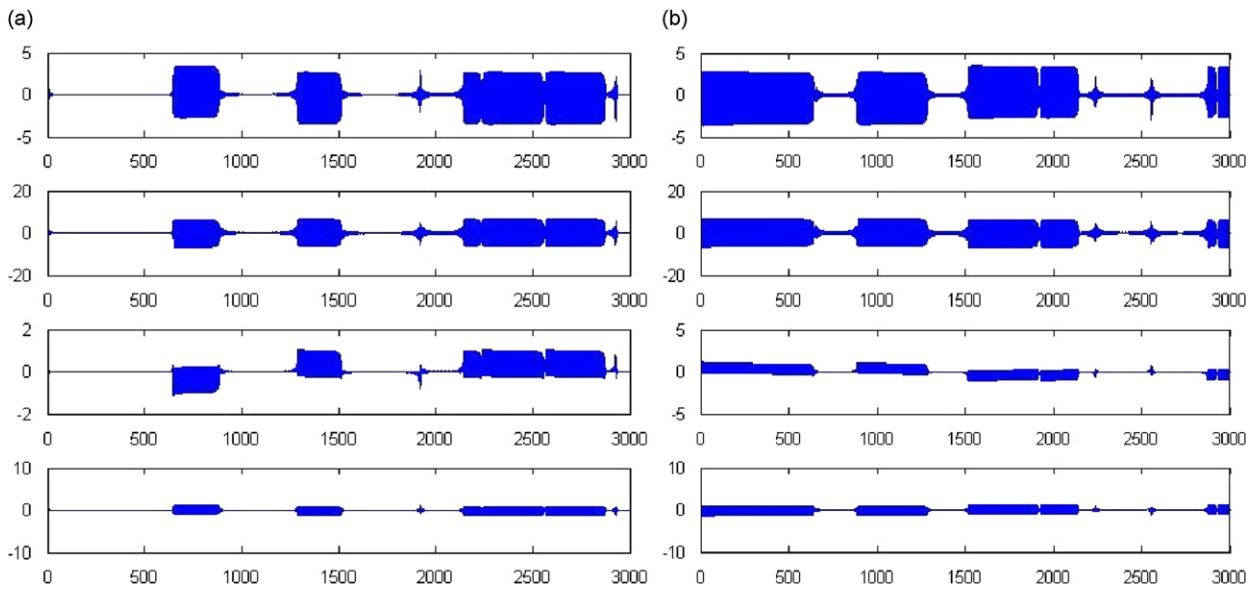


Fig. 5. Alternative CS and AS for initial condition  $(x_1, y_1, u_1, v_1) = (2, 5, 1, 0.3)$ ,  $(x_2, y_2, u_2, v_2) = (-3, 5, 2, 9)$ , and  $p = 12, q = 12$ . (a)  $e_1, e_2, e_3, e_4$  and (b)  $E_1, E_2, E_3, E_4$ .

In order to know how this phenomenon depends upon the initial conditions, different initial conditions are given for fixed driving strength. The results are shown in Figs. 3 and 4. Fig. 3(b) shows that  $E_1, E_2, E_3$ , and  $E_4$  tend to zero. As shown in Fig. 3(a), while the  $e_1, e_2, e_3$ , and  $e_4$  do not go to zero. Comparing Fig. 1 with Fig. 3, it is found that they have contrary behavior. The only reason lies in the different initial conditions. Similar result also exists by comparing Fig. 2 with Fig. 4.

Besides, we also discover the alternative CS and AS. In Fig. 5, the system shows alternative switching between these two states where the initial condition  $(x_1, y_1, u_1, v_1) = (2, 5, 1, 0.3)$ ,  $(x_2, y_2, u_2, v_2) = (-8, -9, 0, 5)$ , and  $p = 12, q = 12$ .

#### 4. Conclusions

In this paper, parameter excited chaos synchronizations of two identical double Duffing systems are studied by adjusting the strength of the substituting variable. Numerical simulations are illustrated for CS or AS of which the occurrence depends on initial conditions and driving strength. Besides, alternative CS and AS is also discovered with the same initial conditions and the same driving strength.

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#### References

- [1] L.M. Pecora, T.L. Carroll, Synchronization in chaotic systems, *Physical Review Letters* 64 (1990) 821–824.
- [2] T.L. Carroll, J.F. Heagy, L.M. Pecora, Transforming signals with chaotic synchronization, *Physical Review E* 54 (1996) 4676–4680.
- [3] L. Kocarev, U. Parlitz, Generalized synchronization, predictability, and equivalence of unidirectionally coupled dynamical systems, *Physical Review Letters* 76 (1996) 1816–1819.
- [4] M.G. Rosenblum, A.S. Pikovsky, J. Kurths, Phase synchronization of chaotic oscillators, *Physical Review Letters* 76 (1996) 1804–1807.
- [5] S.S. Yang, C.K. Duan, Generalized synchronization in chaotic systems, *Chaos, Solitons and Fractals* 9 (1998) 1703–1707.
- [6] G. Chen, S.T. Liu, On generalized synchronization of spatial chaos, *Chaos, Solitons and Fractals* 15 (2003) 311–318.
- [7] C.M. Kim, S. Rim, W.H. Kye, J.W. Ryu, Y.J. Park, Anti-synchronization of chaotic oscillators, *Physics Letters A* 320 (2003) 39–46.
- [8] S.P. Yang, H.Y. Niu, G. Tian, et al., Synchronizing chaos by driving parameter, *Acta Physica Sinica* 50 (2001) 619–623.
- [9] D. Dai, X.K. Ma, Chaos synchronization by using intermittent parametric adaptive control method, *Physics Letters A* 288 (2001) 23–28.
- [10] H.K. Chen, Synchronization of two different chaotic systems: a new system and each of the dynamical systems Lorenz, Chen and Lü, *Chaos, Solitons and Fractals* 25 (2005) 1049–1056.
- [11] H.K. Chen, T.N. Lin, Synchronization of chaotic symmetric gyros by one-way coupling conditions, *ImechE Part C: Journal of Mechanical Engineering Science* 217 (2003) 331–340.
- [12] H.K. Chen, Chaos and chaos synchronization of a symmetric gyro with linear-plus-cubic damping, *Journal of Sound and Vibration* 255 (2002) 719–740.
- [13] Z.M. Ge, T.C. Yu, Y.S. Chen, Chaos synchronization of a horizontal platform system, *Journal of Sound and Vibration* (2003) 731–749.
- [14] Z.M. Ge, T.N. Lin, Chaos, chaos control and synchronization of electro-mechanical gyrostator system, *Journal of Sound and Vibration* 259 (3) (2003).
- [15] Z.M. Ge, Y.S. Chen, Synchronization of unidirectional coupled chaotic systems via partial stability, *Chaos, Solitons and Fractals* 21 (2004) 101–111.
- [16] Z.M. Ge, C.C. Chen, Phase synchronization of coupled chaotic multiple time scales systems, *Chaos, Solitons and Fractals* 20 (2004) 639–647.
- [17] Z.M. Ge, C.C. Lin, Y.S. Chen, Chaos, chaos control and synchronization of vibrometer system, *Journal of Mechanical Engineering Science* 218 (2004) 1001–1020.
- [18] H.K. Chen, T.N. Lin, J.H. Chen, The stability of chaos synchronization of the Japanese attractors and its application, *Japanese Journal of Applied Physics* 42 (12) (2003) 7603–7610.
- [19] Z.M. Ge, Shiue, Non-linear dynamics and control of chaos for Tachometer, *Journal of Sound and Vibration* 253 (4) (2002).
- [20] Z.M. Ge, C.I. Lee, Non-linear dynamics and control of chaos for a rotational machine with a hexagonal centrifugal governor with a spring, *Journal of Sound and Vibration* 262 (2003) 845–864.
- [21] Z.M. Ge, C.M. Hsiao, Y.S. Chen, Non-linear dynamics and chaos control for a time delay Duffing system, *International Journal of Nonlinear Sciences and Numerical Simulation* 6 (2) (2005) 187–199.
- [22] Z.M. Ge, P.C. Tzen, S.C. Lee, Parametric analysis and fractal-like basins of attraction by modified interpolates cell mapping, *Journal of Sound and Vibration* 253 (3) (2002).
- [23] Z.M. Ge, S.C. Lee, Parameter used and accuracies obtain in MICM global analyses, *Journal of Sound and Vibration* 272 (2004) 1079–1085.
- [24] Z.M. Ge, W.Y. Leu, Chaos synchronization and parameter identification for loudspeaker system, *Chaos, Solitons and Fractals* 21 (2004) 1231–1247.
- [25] Z.M. Ge, C.M. Chang, Chaos synchronization and parameter identification for single time scale brushless DC motor, *Chaos, Solitons and Fractals* 20 (2004) 889–903.
- [26] Z.M. Ge, J.K. Lee, Chaos synchronization and parameter identification for gyroscope system, *Applied Mathematics and Computation* 63 (2004) 667–682.

- [27] Z.M. Ge, J.W. Cheng, Chaos synchronization and parameter identification of three time scales brushless DC motor, *Chaos, Solitons and Fractals* 24 (2005) 597–616.
- [28] Z.M. Ge, Y.S. Chen, Adaptive synchronization of unidirectional and mutual coupled chaotic systems, *Chaos, Solitons and Fractals* 26 (2005) 881–888.
- [29] H.K. Chen, Global chaos synchronization of new chaotic systems via nonlinear control, *Chaos, Solitons and Fractals* 4 (2005) 1245–1251.
- [30] H.K. Chen, C.I. Lee, Anti-control of chaos in rigid body motion, *Chaos, Solitons and Fractals* 21 (2004) 957–965.
- [31] Z.M. Ge, H.W. Wu, Chaos synchronization and chaos anticontrol of a suspended track with moving loads, *Journal of Sound and Vibration* 270 (2004) 685–712.
- [32] Z.M. Ge, C.Y. Yu, Y.S. Chen, Chaos synchronization and chaos anticontrol of a rotational supported simple pendulum, *JSME International Journal, Series C* 47 (1) (2004) 233–241.
- [33] Z.M. Ge, W.Y. Leu, Anti-control of chaos of two-degree-of-freedom louder speaker system and chaos system of different order system, *Chaos, Solitons and Fractals* 20 (2004) 503–521.
- [34] Z.M. Ge, J.W. Cheng, Y.S. Chen, Chaos anticontrol and synchronization of three time scales brushless DC motor system, *Chaos, Solitons and Fractals* 22 (2004) 1165–1182.
- [35] Z.M. Ge, C.I. Lee, Anticontrol and synchronization of chaos for an autonomous rotational machine system with a hexagonal centrifugal governor, *Chaos, Solitons and Fractals* 282 (2005) 635–648.
- [36] Z.M. Ge, C.I. Lee, Control, anticontrol and synchronization of chaos for an autonomous rotational machine system with time-delay, *Chaos, Solitons and Fractals* 23 (2005) 1855–1864.
- [37] Jumarie Guy, Fractional master equation: non-standard analysis and Liouville–Riemann derivative., *Chaos, Solitons and Fractals* 12 (2001) 2577–2587.
- [38] H.H. Sun, A.A. Abdelwahad, B. Onaral, *IEEE Transactions on Automatic Control* 29 (1984) 441.
- [39] M. Ichise, Y. Nagayanagi, T. Kojima, *Journal of Electroanalytical Chemistry* 33 (1971) 253.
- [40] O. Heaviside, *Electromagnetic Theory*, Chelsea, New York, 1971.
- [41] P. Arena, R. Caponetto, L. Fortuna, D. Porto, Bifurcation and chaos in noninteger order cellular neural networks, *International Journal of Bifurcation and Chaos* 7 (1998) 1527–1539.
- [42] P. Arena, L. Fortuna, D. Porto, Chaotic behavior in noninteger-order cellular neural networks, *Physical Review E* 61 (2000) 776–781.